Физика на елементарните частици

Lecture 1
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Що е то?

L. Litov Particle Physics Sofia, April 2006



Физика на елементарните частици



Опитва се да отоговори на на два фундаментални въпроса

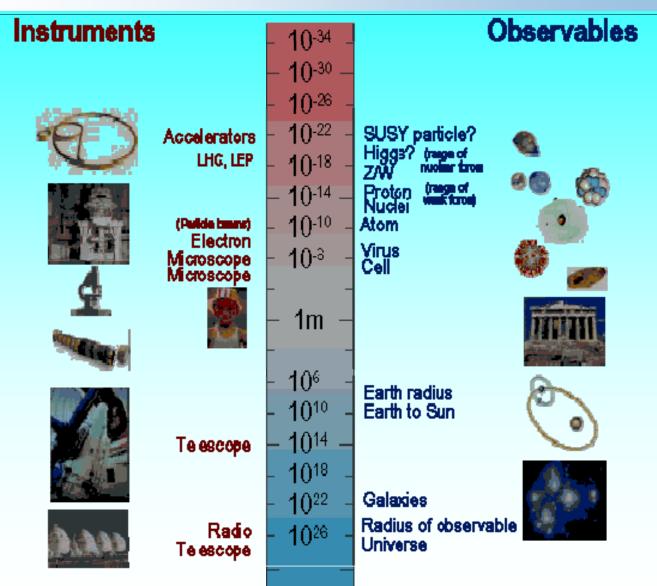
-Кои са елементарните съставящи на материята?

-Кои са фундаменталните сили контролиращи тяхното поведение ?



Мащаби

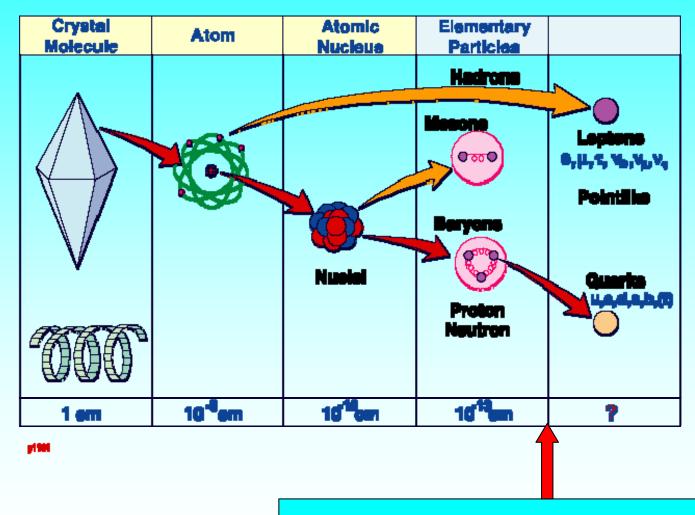






Структура на материята



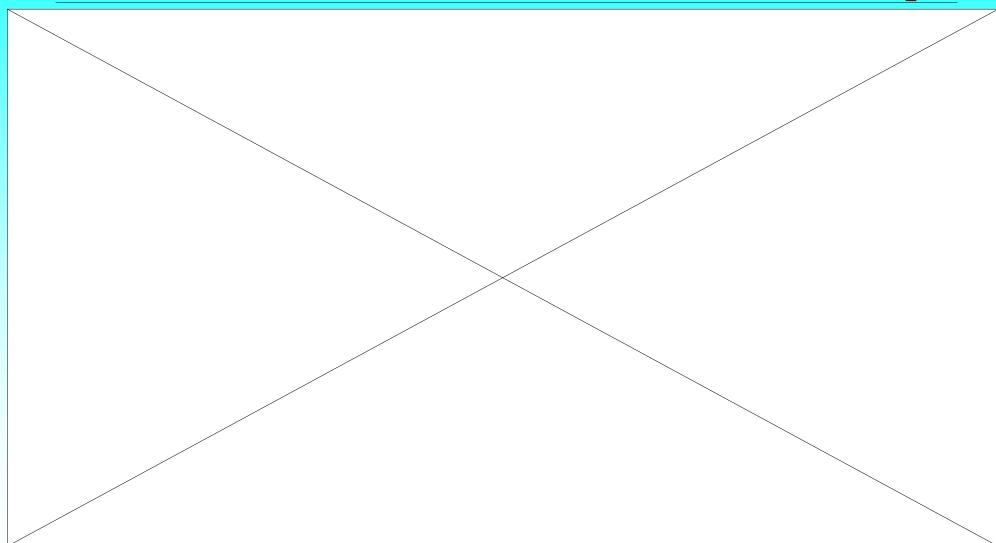


Физика на елементарните частици



Структура на материята







Единици във физиката на елементарните частици

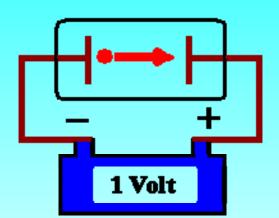


Energy

1 electron-Volt (eV):

the energy of a particle with electric charge = |e|, initially at rest, after acceleration by a difference of electrostatic potential = 1 Volt $(e = 1.60 \times 10^{-19} \text{ C})$

$$1 \text{ eV} = 1.60 \text{ x } 10^{-19} \text{ J}$$



Multiples:

$$1 \text{ keV} = 10^3 \text{ eV}$$
; $1 \text{ MeV} = 10^6 \text{ eV}$

$$1 \text{ GeV} = 10^9 \text{ eV}; \qquad 1 \text{ TeV} = 10^{12} \text{ eV}$$

Energy of a proton in the LHC (in the year 2007):

$$7 \text{ TeV} = 1.12 \times 10^{-6} \text{ J}$$

(the same energy of a body of mass = 1 mg moving at speed = 1.5 m/s)



Маса, разстояние, енергия, температура



These are related quantities

In particle physics the unit of energy is the electron volt.

1 electron volt (eV) = Energy gained by an electron in passing through a voltage difference of 1 V

E=mc² c=3.10⁸ m/s speed of light

E=kT k=10⁻⁴ eV K⁻¹ Boltzmann's constant

E=hc/λ h=4.10-15 eV s Planck's constant

Mass of electron 0.5 million eV (MeV)

Mass of proton 1 Giga eV (GeV)

1 eV ~ 10,000 K 1 GeV ~ 1 femtometre (fm) = 10⁻¹⁵m

General Relativity depends on c and G (Newton's constant), QM depends on ħ. Natural unit of length is given by is called Planck length ~ 10⁻³⁵ m





Фундаментални съставящи

L. Litov Particle Physics Sofia, April 2006



1894 — 1897: <u>Откритие на електрона</u>

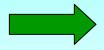
Study of "cathode rays": electric current in tubes at very low gas pressure ("glow discharge")

Measurement of the electron mass: $m_e \approx M_H/1836$

"Could anything at first sight seem more impractical than a body which is so small that its mass is an insignificant fraction of the mass of an atom of hydrogen?" (J.J. Thomson)



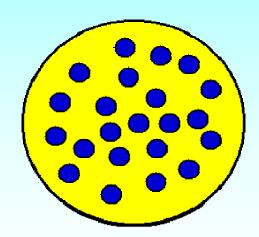
J.J. Thomson



ATOMS ARE NOT ELEMENTARY

Thomson's atomic model:

- Electrically charged sphere
- Radius ~ 10⁻⁸ cm
- Positive electric charge
- Electrons with negative electric charge embedded in the sphere

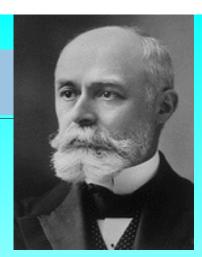




1896: Естествена радиоактивност

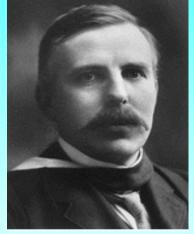
1909 – 13: Rutherford's scattering experiments

Discovery of the atomic nucleus

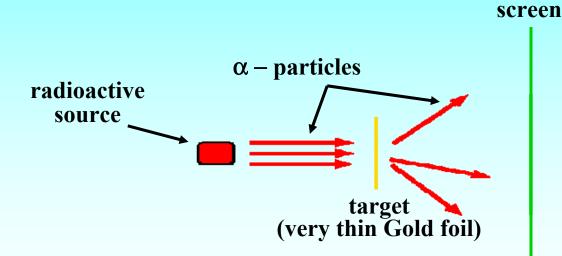


Henri Becquerel

fluorescent



Ernest Rutherford





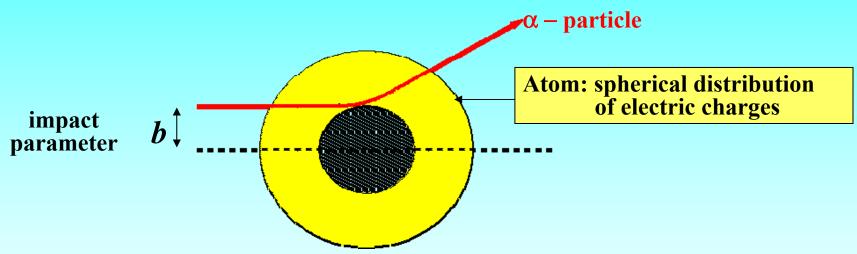
 α – particles: nuclei of Helium atoms spontaneously emitted by heavy radioactive isotopes Typical α – particle velocity $\approx 0.05 \, c$ (c: speed of light)



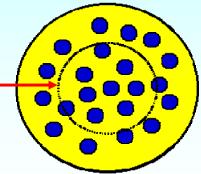
Опити на Ръдърфорд - очаквания



α – atom scattering at low energies is dominated by Coulomb interaction



For Thomson's atomic model the electric charge "seen" by the α – particle is zero, independent of impact parameter



⇒ no significant scattering at large angles is expected

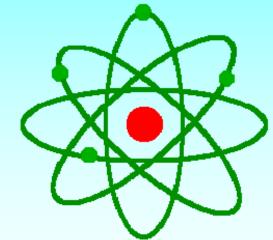


Опити на Ръдърфорд - резултати



significant scattering of α – particles at large angles, consistent with scattering expected for a sphere of radius \approx few x 10⁻¹³ cm and electric charge = $\mathbf{Z}\mathbf{e}$, with $\mathbf{Z}=79$ (atomic number of gold) and $\mathbf{e}=|\text{charge of the electron}|$

an atom consists of a positively charged nucleus surrounded by a cloud of electrons



Nuclear radius $\approx 10^{-13}$ cm $\approx 10^{-5}$ x atomic radius Mass of the nucleus \approx mass of the atom (to a fraction of 1%)



Откритие на неутрона



Neutron: a particle with mass ≈ proton mass but with zero electric charge (Chadwick, 1932) Solution to the nuclear structure problem: Nucleus with atomic number Z and mass number A: a bound system of Z protons and (A – Z) neutrons

Nitrogen anomaly: no problem if neutron spin = $\frac{1}{2}\hbar$ Nitrogen nucleus (A = 14, Z = 7): 7 protons, 7 neutrons = 14 spin $\frac{1}{2}$ particles James Chadwick \Rightarrow total spin has integer value

Neutron source in Chadwick's experiments: a 210 Po radioactive source (5 MeV α – particles) mixed with Beryllium powder \Rightarrow emission of electrically neutral radiation capable of traversing several centimetres of Pb:

4
He₂ + 9 Be₄ → 12 C₆ + neutron
 $^{\alpha}$ - particle



Принцип на Паули



In Quantum Mechanics the electron orbits around the nucleus are "quantized": only some specific orbits (characterized by integer quantum numbers) are possible.

$$R_n = \frac{4\pi\epsilon_0 \hbar^2 n^2}{me^2} \approx 0.53 \times 10^{-10} n^2 \text{ [m]}$$

$$E_n = -\frac{me^4}{2(4\pi\epsilon_0)^2\hbar^2n^2} \approx -\frac{13.6}{n^2} [\text{eV}]$$

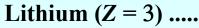
$$m = m_e m_p / (m_e + m_p)$$

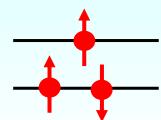
 $n = 1, 2,$

In atoms with Z > 2 only two electrons are found in the innermost orbit – WHY?

ANSWER (Pauli, 1925): two electrons (spin = $\frac{1}{2}$) can never be in the same physical state

Hydrogen
$$(Z = 1)$$
 Helium $(Z = 2)$ Lithium $(Z = 3)$







Wolfgang Pauli

Lowest energy state

Pauli's exclusion principle applies to all particles with half-integer spin (collectively named Fermions) Particle Physics



<u>Антиматерия</u>

Discovered "theoretically" by P.A.M. Dirac (1928)

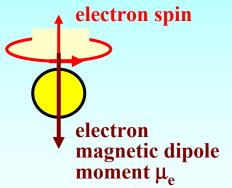
<u>Dirac's equation</u>: a relativistic wave equation for the electron



P.A.M. Dirac

Two surprising results:

■ Motion of an electron in an electromagnetic field:
 presence of a term describing (for slow electrons) the
 potential energy of a magnetic dipole moment in a magnetic field
 ⇒ existence of an intrinsic electron magnetic dipole moment opposite to spin



$$\mu_e = \frac{e\hbar}{2m_e} \approx 5.79 \times 10^{-5} \text{ [eV/T]}$$

■ For each solution of Dirac's equation with electron energy E > 0 there is another solution with E < 0

What is the physical meaning of these "negative energy" solutions?



Експериментално наблюдение на антиматерия

антиматерия

(C.D. Anderson, 1932)

Measure particle momentum and sign of electric charge from magnetic curvature

Carl D. Anderson

<u>Detector</u>: a Wilson cloud – chamber (visual detector based on a gas volume containing vapour close to saturation) in a magnetic field, exposed to cosmic rays

projection of the par

Lorentz force $\vec{f} = e\vec{\mathbf{v}} \times \vec{B}$

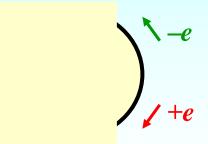
 \rightarrow projection of the particle trajectory in a plane perpendicular to B is a circle

Circle radius for electric charge
$$|e|$$
:

$$p_{\perp}$$
: momentum component perpendicular to magnetic field direction

$$R [m] = \frac{10 p_{\perp} [\text{GeV/c}]}{3B [\text{T}]}$$

NOTE: impossible to distinguish between positively and negatively charged particles going in opposite directions

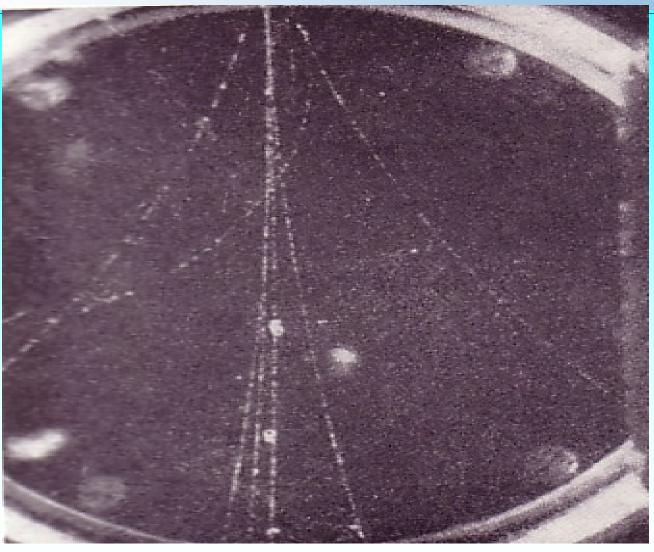


⇒ need an independent determination of the particle direction of motion



Experimental confirmation of antimatter





Cosmic-ray "shower" containing several e e pairs



Неутрино



December 1930: public letter sent by W. Pauli to a physics meeting in Tübingen

Zürich, Dec. 4, 1930

Dear Radioactive Ladies and Gentlemen,

...because of the "wrong" statistics of the N and 6 Li nuclei and the continuous β -spectrum, I have hit upon a desperate remedy to save the law of conservation of energy. Namely, the possibility that there could exist in the nuclei electrically neutral particles, that I wish to call neutrons, which have spin ½ and obey the exclusion principle The mass of the neutrons should be of the same order of magnitude as the electron mass and in any event not larger than 0.01 proton masses. The continuous β -spectrum would then become understandable by the assumption that in β -decay a neutron is emitted in addition to the electron such that the sum of the energies of the neutron and electron is constant. For the moment, however, I do not dare to publish anything on this idea So, dear Radioactives, examine and judge it. Unfortunately I cannot appear in Tübingen personally, since I am indispensable here in Zürich because of a ball on the night of 6/7 December.

W. Pauli



Теория на β-разпада (V-A модел)

 β^- decay: $n \rightarrow p + e^- + \overline{\nu}$

β+ decay: $p \rightarrow n + e^+ + ν$ (e.g., ${}^{14}O_8 \rightarrow {}^{14}N_7 + e^+ + ν$)

v: the particle proposed by Pauli

(named "neutrino" by Fermi)

 \overline{v} : its antiparticle (antineutrino)



Enrico Fermi

(E. Fermi, 1932-33)

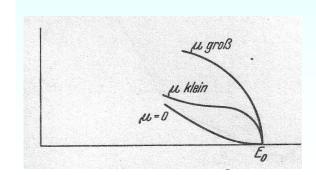
Fermi's theory: a point interaction among four spin ½ particles, using the mathematical formalism of creation and annihilation operators invented by Jordan

 \Rightarrow particles emitted in β – decay need not exist before emission – they are "created" at the instant of decay

Prediction of β – decay rates and electron energy spectra as a function of only one parameter: Fermi coupling constant G_F (determined from experiments)

Energy spectrum dependence on neutrino mass µ (from Fermi's original article, published in German on Zeitschrift für Physik, following rejection of the English version by Nature)

Measurable distortions for $\mu > 0$ near the end-point $(E_0: \max \text{ allowed electron energy})$



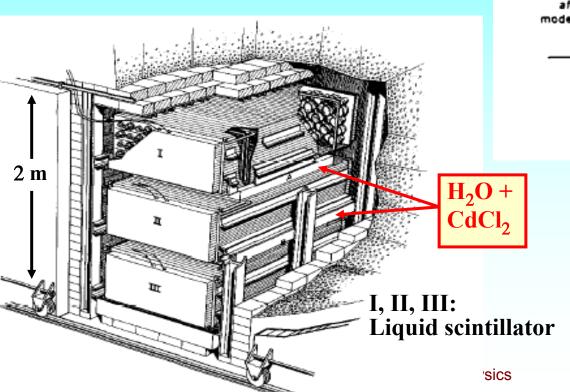


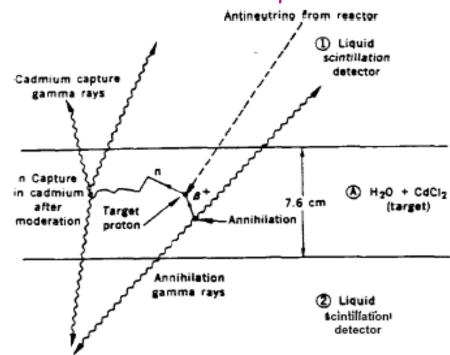
Наблюдение на неутриното



$$\overline{\nu} + p \rightarrow e^+ + n$$

- detect 0.5 MeV γ -rays from $e^+e^- \rightarrow \gamma \gamma$ (t = 0)
- neutron "thermalization" followed by capture in Cd nuclei ⇒ emission of delayed γ-rays (average delay ~30 µs)





(Reines, Cowan 1953)

Event rate at the Savannah River nuclear power plant:

 3.0 ± 0.2 events / hour

(after subracting event rate measured with reactor OFF)

in agreement with expectations
Sofia, April 2006



1947: Откритие на π -мезон



Observation of the $\pi^+ \to \mu^+ \to e^+$ decay chain in a exposed to cosmic rays at high altitudes

In all events the muon has a fixed kinetic energy (4.1 MeV, corresponding to a range of $\sim 600 \, \mu m$ in nuclear emulsion) \Rightarrow two-body decay

 $m_{\pi} = 139.57 \text{ MeV/}c^2$; spin = 0 Dominant decay mode: $\pi^+ \to \mu^+ + \nu$

(and $\pi^- \rightarrow \mu^- + \nu \overline{)}$

Mean life at rest: $\tau_{\pi} = 2.6 \text{ x } 10^{-8} \text{ s} = 26 \text{ ns}$

A neutral π – meson (π °) also exists:

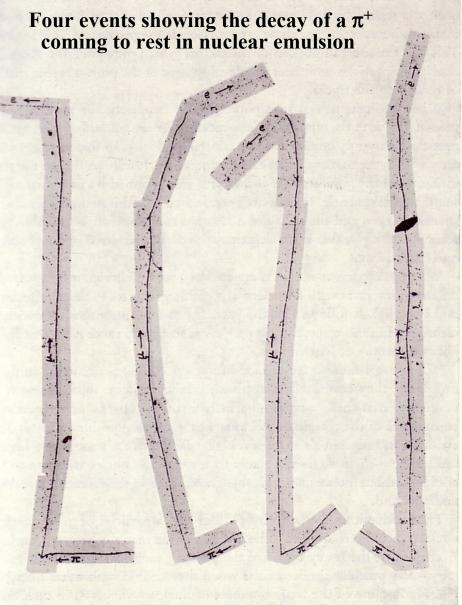
 $m (\pi^{\circ}) = 134.98 \text{ MeV} / c^2$

Decay: $\pi^{\circ} \rightarrow \gamma + \gamma$, mean life = 8.4 x 10^{-17} s

 π – mesons are the most copiously produced particles in proton – proton and proton – nucleus collisions at high energies

L. Litov

Particle Physics





Запазващи се квантови числа



Possible proton decay modes (allowed by all known conservation laws: energy – momentum, electric charge, angular momentum):

$$p \rightarrow \pi^{\circ} + e^{+}$$
 $p \rightarrow \pi^{\circ} + \mu^{+}$
 $p \rightarrow \pi^{+} + \nu$
Why is the free proton stable?

No proton decay ever observed – the proton is STABLE

Limit on the proton mean life: $\tau_p > 1.6 \times 10^{32}$ years

Invent a new quantum number: "Baryonic Number" B

B = 1 for proton, neutron

B = -1 for antiproton, antineutron

B = 0 for e^{\pm} , μ^{\pm} , neutrinos, mesons, photons

Require conservation of baryonic number in all particle processes:

$$\sum_{i} \mathbf{B}_{i} = \sum_{f} \mathbf{B}_{f}$$

(i: initial state particle; f: final state particle)



Странност



Late 1940's: discovery of a variety of heavier mesons (K – mesons) and baryons ("hyperons") – studied in detail in the 1950's at the new high-energy proton synchrotrons (the 3 GeV "cosmotron" at the Brookhaven National Lab and the 6 GeV Bevatron at Berkeley)

Mass values

```
Mesons (spin = 0): m(K^{\pm}) = 493.68 \text{ MeV}/c^2 ; m(K^{\circ}) = 497.67 \text{ MeV}/c^2
Hyperons (spin = ½): m(\Lambda) = 1115.7 \text{ MeV}/c^2 ; m(\Sigma^{\pm}) = 1189.4 \text{ MeV}/c^2
m(\Xi^{\circ}) = 1314.8 \text{ MeV}/c^2; m(\Xi^{-}) = 1321.3 \text{ MeV}/c^2
```

Properties

- Abundant production in proton nucleus, π nucleus collisions
- Production cross-section typical of strong interactions ($\sigma > 10^{-27} \text{ cm}^2$)
- Production in pairs (example: $\pi^- + p \rightarrow K^\circ + \Lambda$; $K^- + p \rightarrow \Xi^- + K^+$)
- Decaying to lighter particles with mean life values $10^{-8} 10^{-10}$ s (as expected for a weak decay)

Examples of decay modes

$$\begin{array}{l} K^{\pm} \to \pi^{\pm} \, \pi^{\circ} \; ; \; K^{\pm} \to \pi^{\pm} \, \pi^{+} \pi^{-} \; ; \; \; K^{\pm} \to \pi^{\pm} \, \pi^{\circ} \, \pi^{\circ} \; ; \; K^{\circ} \to \pi^{+} \pi^{-} \; ; \; K^{\circ} \to \pi^{\circ} \, \pi^{\circ} \; ; \; \ldots \\ \Lambda \to p \; \pi^{-} \; ; \; \; \Lambda \to n \; \pi^{\circ} \; ; \; \Sigma^{+} \to p \; \pi^{\circ} \; ; \; \Sigma^{+} \to n \; \pi^{+} \; ; \; \Sigma^{+} \to n \; \pi^{-} \; ; \; \ldots \\ \Xi \, ^{-} \to \Lambda \, \pi^{-} \; \; ; \; \Xi^{\circ} \to \Lambda \, \pi^{\circ} \end{array}$$



Странност



(Gell-Mann, Nakano, Nishijima, 1953)

• not conserved in weak decays: $S_i - \sum_f S_f = 1$

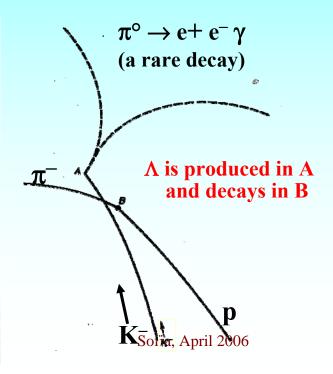
$$\sum_{i} \mathbf{S}_{i} = \sum_{f} \mathbf{S}_{f}$$

S = +1: K^+ , K° ; S = -1: Λ , Σ^{\pm} , Σ° ; S = -2: Ξ° , Ξ^- ; S = 0: all other particles (and opposite strangeness –S for the corresponding antiparticles)

Example of a K - stopping in liquid hydrogen:

$$K^- + p \rightarrow \Lambda + \pi^{\circ}$$
 (strangeness conserving) followed by the decay

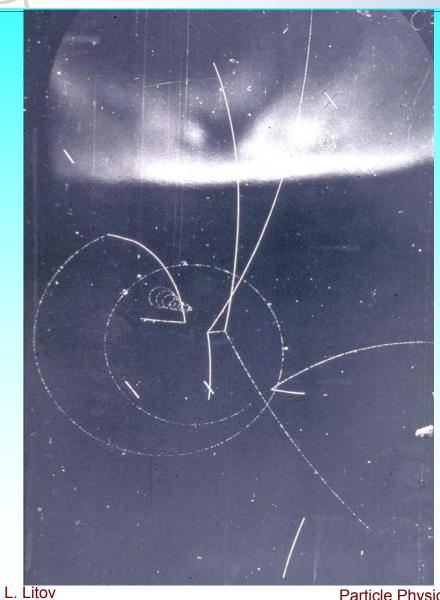
$$\Lambda \rightarrow p + \pi^-$$
 (strangeness violation)

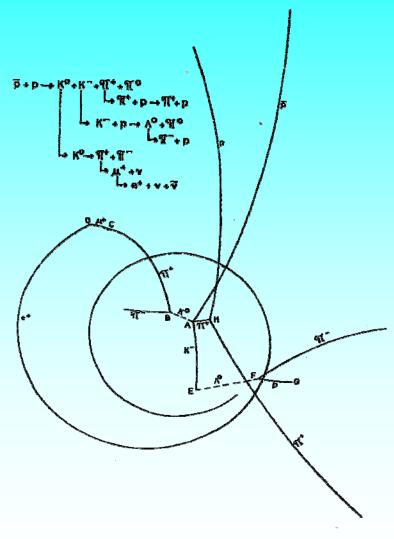




Откритие на антипротона









Лептонно число



A puzzle of the late 1950's: the absence of $\mu \rightarrow e \gamma$ decays

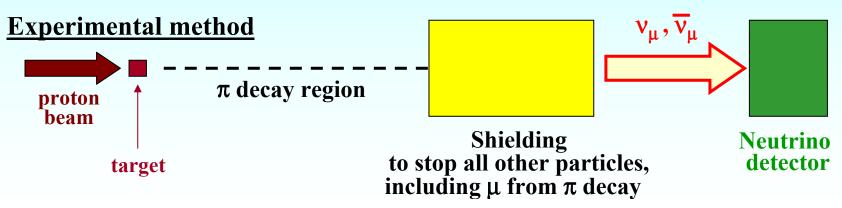
Experimental limit: < 1 in 10^6 $\mu^+ \rightarrow e^+ \nu \overline{\nu}$ decays

A possible solution: existence of a new, conserved "muonic" quantum number distinguishing muons from electrons

To allow $\mu^+ \rightarrow e^+ \nu \ \overline{\nu}$ decays, $\overline{\nu}$ must have "muonic" quantum number but not $v \Rightarrow in \mu^+$ decay the \overline{v} is not the antiparticle of v

 $\Rightarrow \text{two distinct neutrinos } (\nu_e \ , \nu_\mu) \text{ in the decay} \qquad \mu^+ \to e^+ \ \nu_e \ \overline{\nu}_\mu$ Consequence for π – meson decays: $\pi^+ \to \mu^+ \ \nu_\mu \ ; \ \pi^- \to \mu^- \overline{\nu}_\mu$ to conserve the "muonic" quantum number

High energy proton accelerators: intense sources of π^{\pm} mesons $\Rightarrow \nu_{\mu}$, $\overline{\nu}_{\mu}$



If $v_{\mu} \neq v_{e}$, v_{μ} interactions produce μ^{-} and not e^{-} (example: $v_{\mu} + n \rightarrow \mu^{-} + p$)



Кварков модел



Late 1950's – early 1960's: discovery of many strongly interacting particles at the high energy proton accelerators (Berkeley Bevatron, BNL AGS, CERN PS), all with very short mean life times $(10^{-20} - 10^{-23})$ s, typical of strong decays) \Rightarrow catalog of > 100 strongly interacting particles (collectively named "hadrons")

ARE HADRONS ELEMENTARY PARTICLES?

1964 (Gell-Mann, Zweig): Hadron classification into "families"; observation that all hadrons could be built from three spin ½ "building blocks" (named "quarks" by Gell-Mann):

	u	d	S
Electric charge (units e)	+2/3	-1/3	-1/3
Baryonic number	1/3	1/3	1/3
Strangeness	0	0	-1

and three antiquarks $(\overline{u}, \overline{d}, \overline{s})$ with opposite electric charge and opposite baryonic numbers and strangeness



Кварков модел



Mesons: quark – antiquark pairs

Examples of non-strange mesons:

$$\pi^+ \equiv u\overline{d}$$
 ; $\pi^- \equiv \overline{u}d$; $\pi^0 \equiv (d\overline{d} - u\overline{u})/\sqrt{2}$

Examples of strange mesons:

$$K^{-} \equiv s\overline{u}$$
 ; $\overline{K}^{0} \equiv s\overline{d}$; $K^{+} \equiv \overline{s}u$; $K^{0} \equiv \overline{s}d$

Baryons: three quarks bound together

Antibaryons: three antiquarks bound together

Examples of non-strange baryons:

proton
$$\equiv uud$$
 ; neutron $\equiv udd$

Examples of strangeness -1 baryons:

$$\Sigma^+ \equiv suu$$
 ; $\Sigma^0 \equiv sud$; $\Sigma^- \equiv sdd$

Examples of strangeness -2 baryons:

$$\Xi^0 \equiv ssu$$
; $\Xi^- \equiv ssd$



A success of the static quark model



The "decuplet" of spin $\frac{3}{2}$ baryons

Strangeness							Mass (MeV/c ²)
0	N* ⁺⁺ uuu	N* ⁺ uud	N*° udd		N* ⁻ ddd	1232	
-1	Σ* ⁺ suu		Σ*° sud		Σ^{*-} sdd		1384
-2		Ξ*°		Ξ*- ssd			1533
-3			Ω^-				1672 (predicted)

 Ω^- : the bound state of three s – quarks with the lowest mass with total angular momentum = $3/2 \implies 0$

Pauli's exclusion principle requires that the three quarks cannot be identical

Particle Physics



A success of the static quark model



The first Ω^- event (observed in the 2 m liquid hydrogen bubble chamber at BNL using a 5 GeV/c K⁻ beam from the 30 GeV AGS)

Chain of events in the picture:

$$K^- + p \rightarrow \Omega^- + K^+ + K^\circ$$
 (strangeness conserving)

$$\Omega^- \rightarrow \Xi^\circ + \pi^-$$

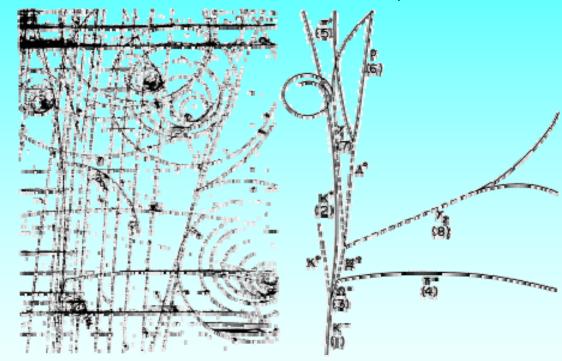
($\Delta S = 1$ weak decay)

$$\Xi^{\circ} \rightarrow \pi^{\circ} + \Lambda$$

($\Delta S = 1$ weak decay)

$$\Lambda \rightarrow \pi^- + p$$

($\Delta S = 1$ weak decay)



$$\pi^{\circ} \rightarrow \gamma + \gamma$$
 (electromagnetic decay)

with both γ – rays converting to an e^+e^- in liquid hydrogen

(very lucky event, because the mean free path for $\gamma \rightarrow e^+e^-$ in liquid hydrogen is ~10 m)

Ω^- mass measured from this event = $1686 \pm 12 \text{ MeV}/c^2$



Дълбоко нееластично разсейване

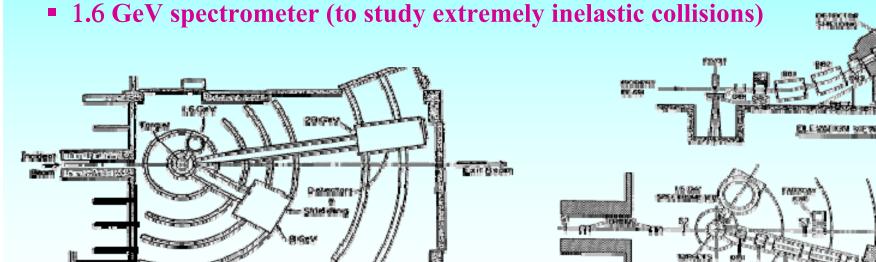


Electron – proton scattering using a 20 GeV electron beam from the Stanford two – mile Linear Accelerator (1968 – 69).

The modern version of Rutherford's original experiment: resolving power \approx wavelength associated with 20 GeV electron $\approx 10^{-15}$ cm

Three magnetic spectrometers to detect the scattered electron:

- 20 GeV spectrometer (to study elastic scattering $e^- + p \rightarrow e^- + p$)
- 8 GeV spectrometer (to study inelastic scattering $e^- + p \rightarrow e^- + hadrons$)





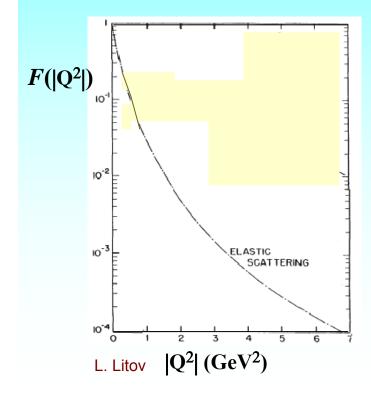
Дълбоко нееластично разсейване



Electron elastic scattering from a point-like charge |e| at high energies: differential cross-section in the collision centre-of-mass (Mott's formula)

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 (\hbar c)^2}{8E^2} \frac{\cos^2(\theta/2)}{\sin^4(\theta/2)} \equiv \sigma_M \qquad \left(\alpha = \frac{e^2}{\hbar c} \approx \frac{1}{137}\right)$$

Scattering from an extended charge distribution: multiply σ_M by a "form factor":



$$\frac{d\sigma}{d\Omega} = F(\left|\mathbf{Q}^2\right|)\sigma_M$$

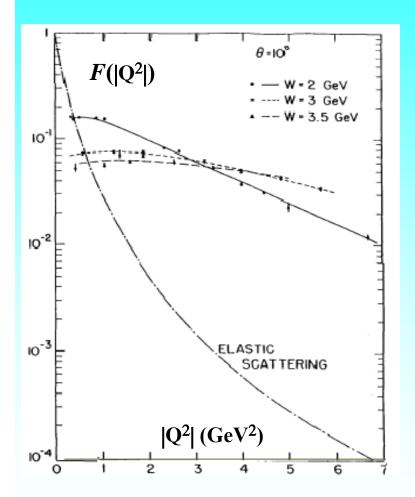
 $|Q| = \hbar / D$: mass of the exchanged virtual photon D: linear size of target region contributing to scattering Increasing $|Q| \Rightarrow$ decreasing target electric charge

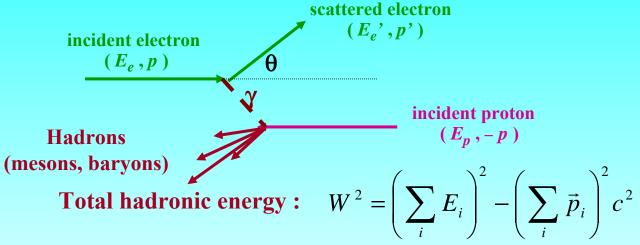
 $F(|\mathbf{Q}^2|) = 1$ for a point-like particle \Rightarrow the proton is not a point-like particle



Дълбоко нееластично разсейване







For deeply inelastic collisions, the cross-section depends only weakly on $|Q^2|$, suggesting a collision with a <u>POINT-LIKE</u> object



Партонен модел



Deep inelastic electron – proton collisions are elastic collisions with point-like, electrically charged, spin $\frac{1}{2}$ constituents of the proton carrying a fraction x of the incident proton momentum

Each constituent type is described by its electric charge e_i (units of |e|) and by its x distribution (dN_i/dx) ("structure function")

If these constituents are the u and d quarks, then deep inelastic e – p collisions provide information on a particular combination of structure functions:

$$\left(\frac{dN}{dx}\right)_{e-p} = e_u^2 \frac{dN_u}{dx} + e_d^2 \frac{dN_d}{dx}$$

Comparison with $v_u - p$ and $\overline{v}_u - p$ deep inelastic collisions at high energies under the assumption that these collisions are also elastic scatterings on quarks $\nu_{\mu} + p \rightarrow \mu^- + hadrons$: $\nu_{\mu} + d \rightarrow \mu^- + u$ (depends on dN_d/dx)

$$v_{\mu} + p \rightarrow \mu^{-} + \text{hadrons}: v_{\mu} + d \rightarrow \mu^{-} + u \text{ (depends on } dN_{d} / dx \text{)}$$

$$\overline{v_{\mu}} + p \rightarrow \mu^{+} + \text{hadrons}: v_{\mu} + u \rightarrow \mu^{+} + d \text{ (depends on } dN_{u} / dx \text{)}$$

(Neutrino interactions do not depend on electric charge)

All experimental results on deep inelastic e - p, $v_u - p$, $v_u - p$ collisions are consistent with $e_u^2 = 4/9$ and $e_d^2 = 1/9$



the proton constitution tysia re the quarks

Sofia, April 2006



Цвят



Problem with

J=3/2 barion
$$\Delta^{++} \Rightarrow u \uparrow u \uparrow u \uparrow$$

It has symmetric wave function, but it is a fermion Contradiction with Pauli exclusion principle the wave function should be antisymmetric The way out - new quantum number - colour

N_c = 3
$$q^{\alpha}$$
 , α = 1,2,3
Then $\Delta^{++} = \frac{1}{\sqrt{6}} \varepsilon^{\alpha\beta\gamma} u_{\alpha} \uparrow u_{\beta} \uparrow u_{\gamma} \uparrow$

$$B = \frac{1}{\sqrt{6}} \varepsilon^{\alpha\beta\gamma} | q_{\alpha} q_{\beta} q_{\gamma} >$$

In general case 1
$$B = \frac{1}{\sqrt{6}} \varepsilon^{\alpha\beta\gamma} | q_{\alpha} q_{\beta} q_{\gamma} > \qquad M = \frac{1}{\sqrt{3}} \delta^{\alpha\beta} | q_{\alpha} \overline{q}_{\beta} >$$



Експериментална проверка на хипотезата за цвят



Two beams circulating in opposite directions in the same magnetic ring and colliding head-on

$$E, \overrightarrow{p}$$
 $E, -p$
 $E, -p$

A two-step process: $e^{+} + e^{-} \rightarrow virtual photon \rightarrow f + \overline{f}$

f: any electrically charged elementary spin $\frac{1}{2}$ particle (μ , quark) (excluding e⁺e⁻ elastic scattering)

Virtual photon energy – momentum : $E_{\gamma} = 2E$, $p_{\gamma} = 0 \Rightarrow Q^2 = E_{\gamma}^2 - p_{\gamma}^2 c^2 = 4E^2$

Cross - section for
$$e^+e^- \rightarrow f f$$
:

$$\alpha = e^2/(\hbar c) \approx 1/137$$

 e_f : electric charge of particle f (units |e|) $\beta = v/c$ of outgoing particle f

$$\vec{\beta} = v/c$$
 of outgoing particle f

Cross - section for
$$e^+e^- \rightarrow f f$$
:
$$\sigma = \frac{2\pi\alpha^2\hbar^2c^2}{3Q^2}e_f^2\beta(3-\beta)$$

$$\alpha = e^2/(\hbar c) \approx 1/137$$

(formula precisely verified for $e^+e^- \rightarrow \mu^+\mu^-$)

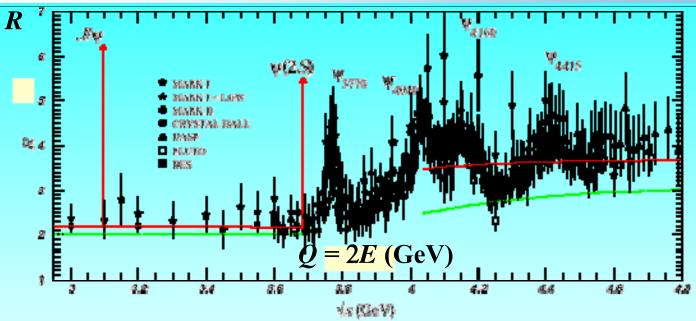
Assumption: $e^+e^- \rightarrow quark(q) + antiquark(\overline{q}) \rightarrow hadrons$ \Rightarrow at energies $E >> m_q c^2$ (for q = u, d, s) $\beta \approx 1$:

$$R = \frac{\sigma(e^{+}e^{-} \to \text{hadrons})}{\sigma(e^{+}e^{-} \to \mu^{+}\mu^{-})} = e_{u}^{2} + e_{d}^{2} + e_{s}^{2} = \frac{4}{9} + \frac{1}{9} + \frac{1}{9} = \frac{2}{3}_{\text{, April 2006}}$$



Stanford e⁺e⁻ collider SPEAR (1974 –75):





- For Q < 3. 6 GeV $R \approx 2$. If each quark exists in three different states, $R \approx 2$ is consistent with 3 x (2/3). This would solve the Ω^- problem.
- Between 3 and 4.5 GeV, the peaks and structures are due to the production of quark-antiquark bound states and resonances of a fourth quark ("charm", c) of electric charge +2/3
- Above 4.6 GeV $R \approx 4.3$. Expect $R \approx 2$ (from u, d, s) + 3 x (4 / 9) = 3.3 from the addition of the c quark alone. So the data suggest pair production of an additional elementary spin ½ particle with electric charge = 1 (later identified as the τ lepton (no strong interaction) with mass ≈ 1777 MeV/ c^2).



Фундаментални частици





